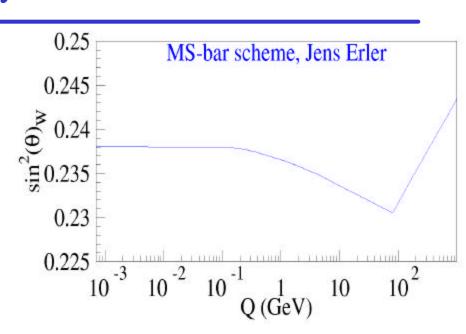
# DIS-Parity: Physics Beyond the Standard Model with Parity NonConserving Deep Inelastic Scattering

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- Introduction: Weinberg-Salam Model and  $sin^2(\theta_w)$
- Parity NonConserving Electron
   Deep Inelastic Scattering
- 11 GeV Measurement at Jefferson Laboratory



Work done in collaboration with Peter Bosted, Dave Mack *et al.* 

# Weinberg-Salam model and $\sin^2(\theta_W)$

Unification of Weak and E&M Force

- •SU(2)—weak isospin—Triplet of gauge bosons
- •U(1)—weak hypercharge—Single gauge boson Electroweak Lagrangian:

 $\mathcal{L} = g \vec{J}_{\mu} \cdot \vec{W}_{\mu} + g' J_{\mu}^{Y} B_{\mu} | J_{\mu}^{Y} = J_{\mu}^{\text{EM}} - J_{\mu}^{(3)}$ 

 $J_m$ ,  $J_m$  isospin and hypercharge currents g, go couplings between currents and fields

$$W_{\mu}^{\pm} = \frac{1}{\sqrt{2}} \left( W_{\mu}^{(1)} \pm i W_{\mu}^{(2)} \right) \qquad \text{Weak CC}$$

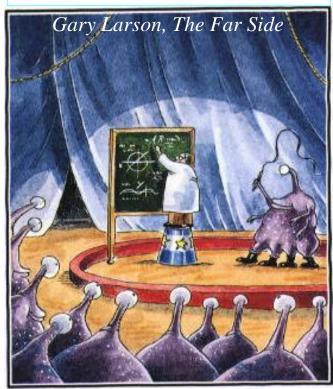
$$A_{\mu} = \frac{1}{\sqrt{g^2 + g'^2}} \left( g' W_{\mu}^{(3)} + g B_{\mu} \right) \qquad \text{EM NC}$$

$$Z_{\mu}^{0} = \frac{1}{\sqrt{g^2 + g'^2}} \left( g' W_{\mu}^{(3)} - g B_{\mu} \right) \qquad \text{Weak NC}$$

 $\theta_{\rm w}$ , relative strength of the SU(2) and U(1) couplings:

$$\tan \theta_W = \frac{g'}{g} \quad \sin \theta_W = \frac{g'}{\sqrt{g'^2 + g^2}}$$
$$\cos \theta_W = \frac{g}{\sqrt{g'^2 + g^2}}$$

Remember—I'm not the expert. . .



Abducted by an alien circus company, Professor Doyle is forced to write calculus equations in center ring.

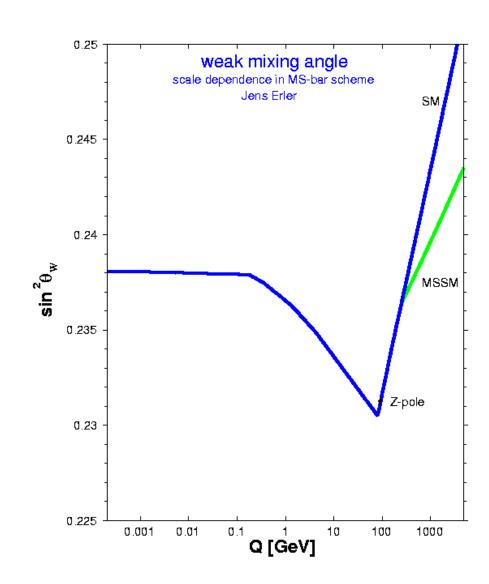
#### •Observables:

•
$$Q_{EM}$$
  $e = g \sin(\theta_W)$ 

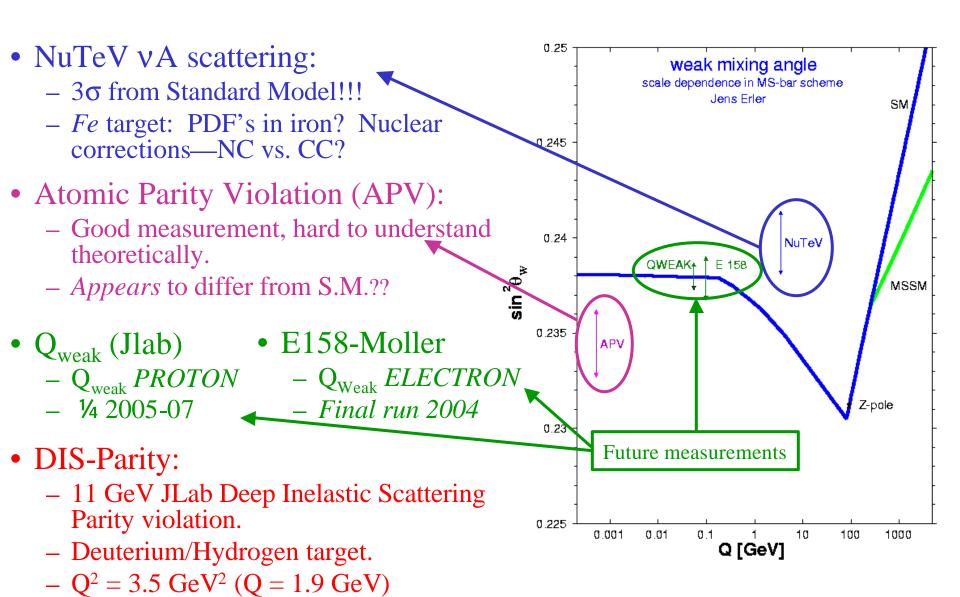
•
$$\sin^2(\theta_W) = 1 - M_W^2 / M_Z^2$$
.

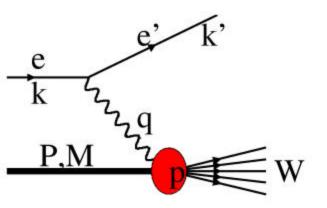
# $\sin^2(\theta_{\rm W})$ vs. $Q^2$

- Standard Model predicts  $\sin^2(\theta_w)$  varies (runs) with  $Q^2$ 
  - Well measured at Z-pole, but not at other Q<sup>2</sup>.
  - Running sensitive to non-Standard Model Physics.
  - Different measurements sensitive to *different* non-S.M. physics.
- $\sin^2(\theta_W)$  is *scheme dependent* observable—it's value depends on the renormalization scheme.



# $\sin^2(\theta_w)$ measurements below Z-pole





$$Q^{2} = -q^{2} = 2(EE^{0} - \mathbf{k}^{0} \mathbf{k})$$

$$-m_{l}^{2} - m_{l}^{2}$$

$$1/4 \ 4EE^{0} \sin^{2}(\theta/2)$$

$$\mathbf{v} = \mathbf{q}^{0} \mathbf{P}/\mathbf{M} = \mathbf{E} - \mathbf{E}^{0}$$

$$x = Q^2/2Mv$$

$$y = q P k P = v / E$$

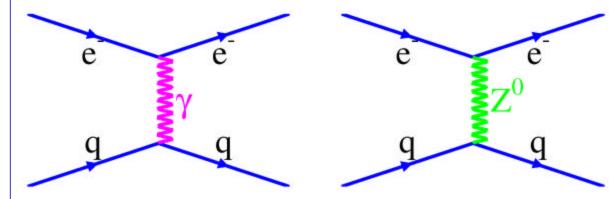
$$W^2 = (\mathbf{P} + \mathbf{q})^2$$
$$= M^2 + 2M\nu - Q^2$$

$$\mathbf{s} = (\mathbf{k} + \mathbf{P})^2$$
$$= \mathbf{Q}^2/\mathbf{x}\mathbf{y} + \mathbf{M}^2 + \mathbf{m}_I^2$$

#### Polarized e deuterium DIS

Look for left-right asymmetry in polarized eD deep inelastic scattering

•Asymmetry caused by interference between  $Z^0$  and  $\gamma$  diagrams.



- •Use deuterium target:  $u(x) \cdot d(x)$
- •Large asymmetry: A<sub>d</sub>¼ 10<sup>-4</sup>

#### **DIS Formalism**

$$A_d = \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R}$$
$$= -\left(\frac{3G_F Q^2}{\sigma_R Q^2}\right)^{\frac{1}{2}}$$

Longitudinally polarized electrons on unpolarized isoscaler (deuterium) target (derivation is problem for listener).

$$= -\left(\frac{3G_FQ^2}{\pi\alpha^2\sqrt{2}}\right)\frac{2C_{1u} - C_{1d}\left[1 + R_s(x)\right] + Y\left(2C_{2u} - C_{2d}\right)R_v(x)}{5 + R_s(x)}$$

$$Y = \frac{1 - (1 - y)^2}{1 + (1 - y)^2 - y^2 R / (1 + R)}$$

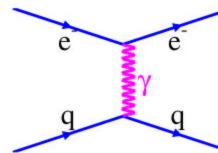
$$R(x,Q^2) = \sigma_L/\sigma_R \approx 0.2$$

$$Y = \frac{1 - (1 - y)^2}{1 + (1 - y)^2 - y^2 R / (1 + R)} \qquad R_s(x) = \frac{2s(x)}{u(x) + d(x)} \xrightarrow{\text{large} \times} 0$$

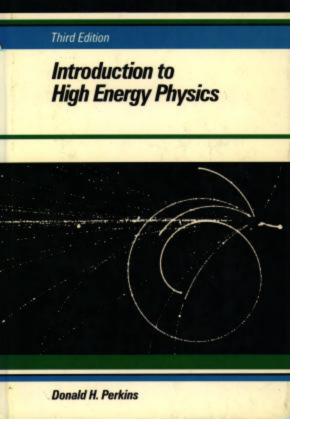
$$R(x, Q^2) = \sigma_L / \sigma_R \approx 0.2 \qquad R_v(x) = \frac{u_v(x) + d_v(x)}{u(x) + d(x)} \xrightarrow{\text{large} \times} 1$$

$$C_{2q}$$
) NC axial coupling to  $q$   
£ NC vector coupling to  $e$ 

$$\begin{split} & C_{1u} = -\frac{1}{2} + \frac{4}{3}\sin^2(\theta_W) \approx -0.19 \\ & C_{1d} = \frac{1}{2} - \frac{2}{3}\sin^2(\theta_W) \approx 0.35 \\ & C_{2u} = -\frac{1}{2} + 2\sin^2(\theta_W) \approx -0.04 \\ & C_{2d} = \frac{1}{2} - 2\sin^2(\theta_W) \approx 0.04. \end{split}$$



Note that each of the  $C_{ia}$  are sensitive to different possible S.M. extensions.



# Textbook Physics: Polarized e<sup>-</sup> d scattering

9.7. Experimental Tests of Neutral Currents in the Weinberg-Salam Model

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#### 9.7.4. Asymmetries in the Scattering of Polarized Electrons by Deuterons

Finally we discuss a very delicate experiment to detect tiny parity-violation effects (asymmetries) due to the interference between  $Z^0$  and  $\gamma$ -exchange in inelastic scattering of polarized electrons by deuterons. The experiment was carried out with beams of electrons of 16-22-GeV/c momentum at SLAC, the reaction being

$$e_{L,R}^- + d_{\text{unpolarized}} \rightarrow e^- + X,$$

# Repeat SLAC experiment (30 years later) with better statistics and systematics at 12 GeV Jefferson Lab:

• Beam current 100 μA vs. 4 μA at SLAC in '78 £ 25

£ 25 stat

• 60 cm target vs. 30 cm target

£ 2 stat

• P<sub>e</sub> (=electron polarization) = 80% vs. 37%

£ 4 stat

• δ P<sub>e</sub> ¼ 1% vs. 6%

£ 6 sys

## **Experimental Constraints and Kinematics**

- Small sea quark uncertainties x > 0.3
- Better sensitivity to  $\sin^2(\theta_{\mathbf{W}})$  ) Large Y
- DIS region, minimize higher twist ) Q<sup>2</sup>>2.0 GeV<sup>2</sup>
  - $W^2>4.0 \text{ GeV}^2$
- d(x)/u(x) uncertainties ) deuterium target
- Pion and other backgrounds )  $E^{0}/E>0.3$  (y<0.7)

Quick calculations show that these conditions are best matched with an 11 GeV beam and an electron scattering angle of approximately 10<sup>±</sup>-15<sup>±</sup> (12.5<sup>±</sup>).

$$hxi = 0.45$$
  $hQ^2i = 3.5 \text{ GeV}^2$ 

$$hYi = 0.46$$
  $hW^2i = 5.23 \text{ GeV}^2$ 

$$\frac{\delta \sin^2 \theta_W}{\sin^2 \theta_W} \bigg|_{Y=0.46} \approx \frac{1}{2} \left( \frac{\delta A_d}{A_d} \right) \qquad A_d \approx 2.9 \times 10^{-4}$$

### Detector and Expected Rates

- Expt. Assumptions:
  - 60 cm ld<sub>2</sub>/lH<sub>2</sub> target
  - 11 GeV beam @ 90μA
  - 75% polar.
  - − 12.5<sup>±</sup> central angle
  - $-12 \, \mathrm{msr} \, \mathrm{d}\Omega$
  - 6.8 GeV§ 10% momentum bite

- Rate expectations:
  - 1MHz DIS
  - $-\pi/e \frac{\pi}{1}$  1 MHz pions
  - 2 MHz Total rate
  - dA/A = 0.5% ) 345 hrs (ideal) plus time for H<sub>2</sub> and systematics studies.

- Will work in either Hall C (HMS +SHMS) or Hall A (MAD)
- $\pi$ /e separation requires gas Cherenkov counters ¼ 6 GeV thresh.
- Ignore tracking in detectors
- Rate requires flash ADC's on Cherenkov and Calorimeters—this is a counting experiment!!

## Uncertainties in A<sub>d</sub>

#### • Beam Polarization:

- QWeak also needs 1.4% polarization accuracy.
- Hall C Moller has achieved
   0.5% polarization accuracy.
- Higher twists may enter in at this low of Q<sup>2</sup>:
  - Check by taking additional data at lower Q<sup>2</sup>
  - 12.5<sup>±</sup>—11 GeV and
     15<sup>±</sup>—8 GeV data
  - Possible 6 GeV experiment?
- EMC effect in d<sub>2</sub>
  - Check with proton data in region where d/u is known.

Statistical	0.5%
Beam polarization	1.0%
$\delta Q^2$	0.5%
Radiative corr.	<1%
$\delta R = \delta(\sigma_L/\sigma_T) = $ \$ 15%	<0.02%
$\delta s(x) = $ § 10%	<0.03%
Higher Twist	????
EMC Effect	????

# Expected $\sin^2(\theta_W)$ Results

$$A = f \left[ \alpha + \beta \sin^2(\theta_W) \right] \quad A = 1.1 \times 10^{-4} Q^2 \left[ 2.2 - 6.1 \sin^2(\theta_W) \right]$$

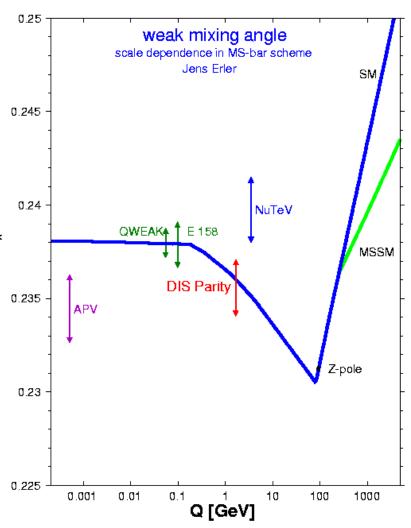
$$\frac{\delta \sin^2 (\theta_W)}{\sin^2 (\theta_W)} = \frac{\delta A}{A} \frac{1}{\beta} \frac{\alpha + \beta \sin^2 (\theta_W)}{\sin^2 (\theta_W)}$$

Measure A<sub>d</sub> to § 0.5% stat § 1.1% syst. (1.24% combined)

Measurement uncertainties driven by polarization uncertainties

$$\frac{\delta \sin^2 \theta_W}{\sin^2 \theta_W} \Big|_{Y=0.46} = 0.56 \left(\frac{\delta A_d}{A_d}\right)$$
$$= 0.7\%$$

What about  $C_{ia}$ 's?



# Extracted Signal—It's all in the binning

$$\frac{A_d}{1.1 \times 10^{-4} Q^2} \approx -\left[ (2C_{1u} - C_{1d}) + Y (2C_{2u} - C_{2d}) \right]$$









PDG:  $C_{1u} = -0.209 \$ 0.041$  highly  $C_{1d} = 0.358 \$ 0.037$  correlated

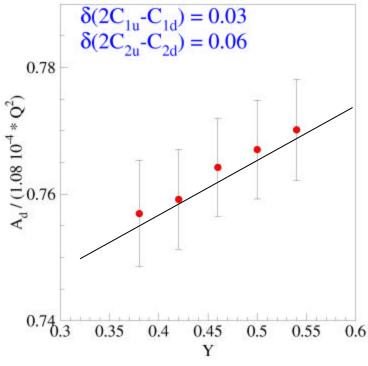
$$2C_{2u} - C_{2d} = -0.08 \$ 0.24$$

#### This measurement:

$$\delta(2C_{10} - C_{1d}) = 0.03$$
 (stat.)

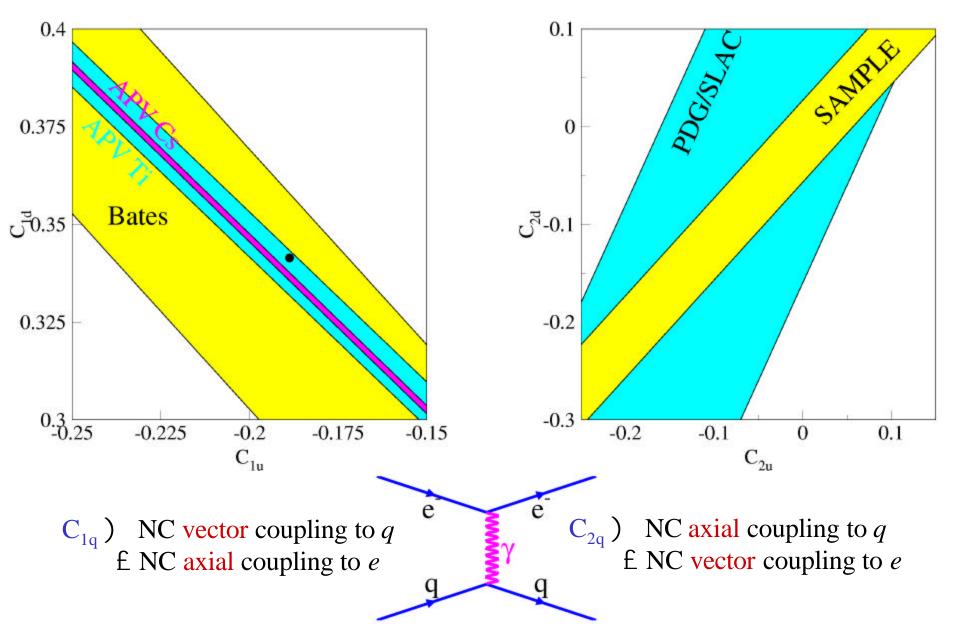
$$\delta(2C_{2u} - C_{2d}) = 0.06$$
 (stat.)

(with out considering other expts.)

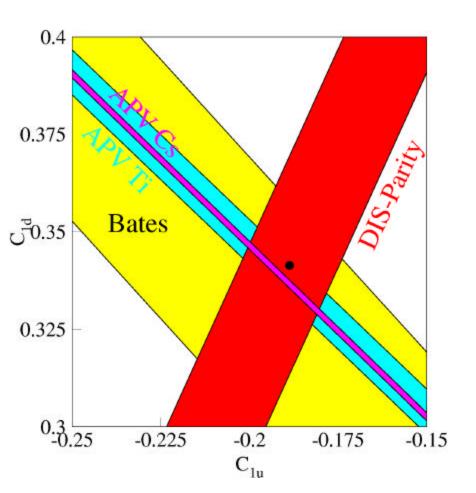


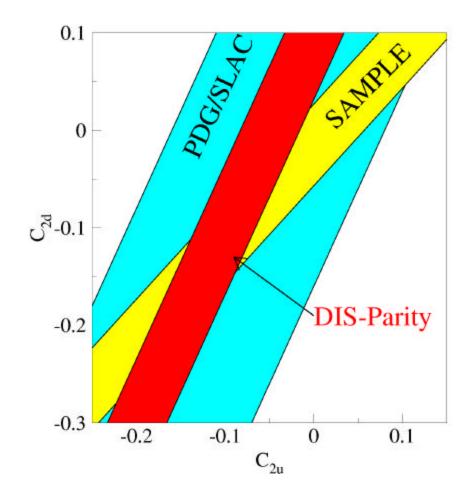
Note—Polarization uncertainty enters as in slope and intercept  $A_{obs} = PA_d / P(2C_{1u}-C_{1d}) + P(2C_{2u}-C_{2d})Y]$ but is correlated

#### Constraints with DIS-Parity



#### Constraints with DIS-Parity

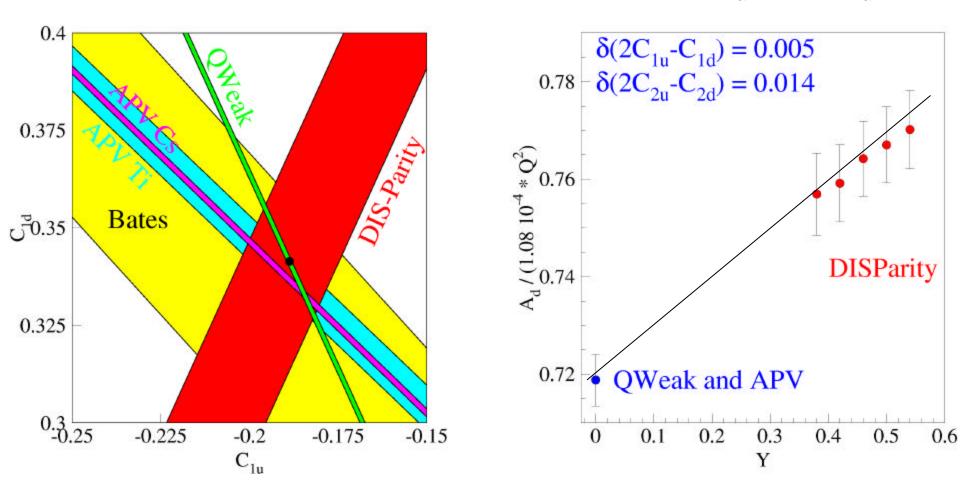




DIS-Parity provides intersecting constraints on C<sub>ia</sub> parameters:

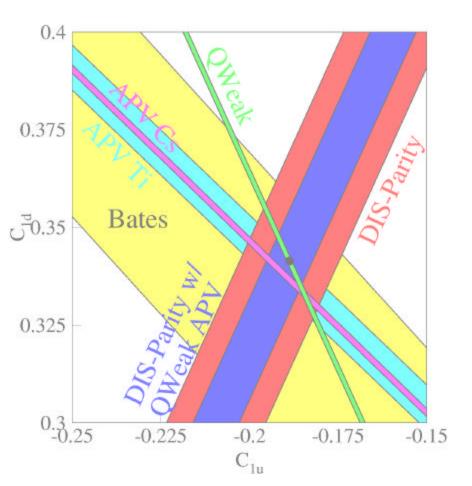
$$\delta(2C_{1u}-C_{1d}) = 0.03 \text{ (stat.)}$$
  $\delta(2C_{2u}-C_{2d}) = 0.06 \text{ (stat.)}$  (1 $\sigma$  limits)

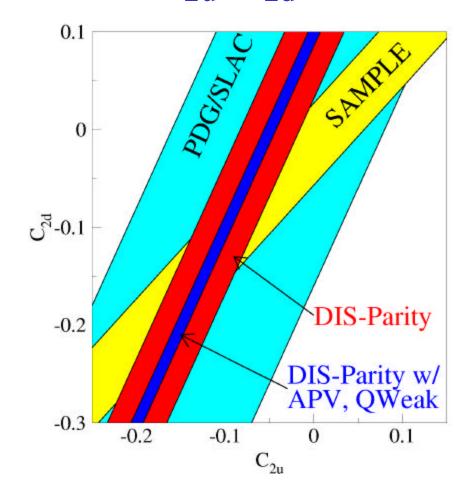
# QWeak & APV will Constrain C<sub>1u</sub> & C<sub>1d</sub>



Combined expected Qweak (proton) and APV measurements give a better value for  $C_{1u}$  and  $C_{1d}$ . Will provide an "anchor" point for fit. Very useful in determining  $2C_{2u} - C_{2d}$ .

# DIS-Parity determines 2C<sub>2u</sub>-C<sub>2d</sub>

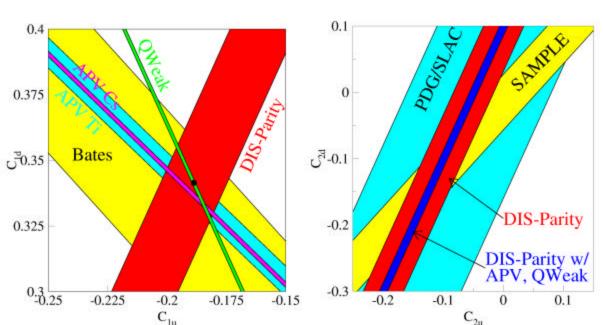


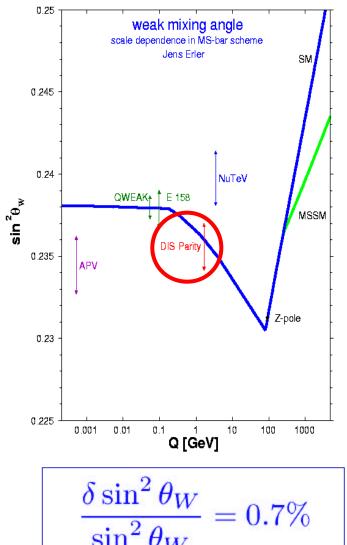


Combined result significantly constrains  $2C_{2u}-C_{2d}$ . PDG  $2C_{2u}-C_{2d}=-0.08$  § 0.24 Combined  $\delta(2C_{2u}-C_{2d})=$  § 0.014 £ 17 improvement (S.M  $2C_{2u}-C_{2d}=0.0986$ )

## **DIS-Parity: Conclusions**

- Measurements of  $\sin^2(\theta_W)$  below  $M_Z$  provide strict tests of the Standard Model.
- Parity NonConserving DIS provides complimentary sensitivity to other planned measurements.
- DIS-Parity Violation measurements can be carried out at Jefferson Lab with the 12 GeV upgrade (beam and detectors) in either Hall A or Hall C.





$$\frac{\delta \sin^2 \theta_W}{\sin^2 \theta_W} = 0.7\%$$
$$\delta(2C_{1u} - C_{1d}) = 0.005$$
$$\delta(2C_{2u} - C_{2d}) = 0.014$$

10 January 2003

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# Weinberg-Salam model and $\sin^2(\theta_W)$

Unification of Weak and E&M Force

- •SU(2)—weak isospin—Triplet of gauge bosons
- •U(1)—weak hypercharge—Single gauge boson

#### Electroweak Lagrangian:

$$\mathcal{L} = g\vec{J}_{\mu} \cdot \vec{W}_{\mu} + g'J_{\mu}^{Y}B_{\mu}$$
$$J_{\mu}^{Y} = J_{\mu}^{\text{EM}} - J_{\mu}^{(3)}$$

 $J_{\mu}^{Y} = J_{\mu}^{\rm EM} - J_{\mu}^{(3)}$  $J_{m} J_{m}^{y}$  isospin and hypercharge currents g, go couplings between currents and fields

$$W_{\mu}^{\pm} = \frac{1}{\sqrt{2}} \left( W_{\mu}^{(1)} \pm i W_{\mu}^{(2)} \right) \quad \text{Weak CC}$$

$$A_{\mu} = \frac{1}{\sqrt{g^2 + g'^2}} \left( g' W_{\mu}^{(3)} + g B_{\mu} \right) \quad \text{EM NC}$$

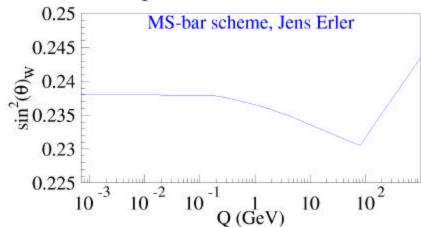
$$Z_{\mu}^{0} = \frac{1}{\sqrt{g^2 + g'^2}} \left( g' W_{\mu}^{(3)} - g B_{\mu} \right) \quad \text{Weak NC}$$

$$\tan \theta_{W} = \frac{g'}{g} \quad \sin \theta_{W} = \frac{g'}{\sqrt{g'^2 + g^2}}$$

$$\cos \theta_{W} = \frac{g}{\sqrt{g'^2 + g^2}}$$

•Observables:  $Q_{EM}$   $e = g \sin(\theta_W)$  $\sin^2(\theta_W) = 1 - M_W^2/M_Z^2$ .

- $\theta_w$ , relative strength of the SU(2) and U(1) couplings:  $\tan(\theta_w)$  g/g
- Standard Model predicts  $\sin^2(\theta_W)$  varies (runs) with  $Q^2$ 
  - Well measured at Z-pole, but not at other Q<sup>2</sup>.



- Running sensitive to non-Standard Model Physics.
- Different measurements sensitive to different non-S.M. physics.
- $\sin^2(\theta_W)$  is *scheme dependent* observable—it's value depends on the renormalization scheme.

# Additional Possibilities with H<sub>2</sub>

- Asymmetry in  $\sigma_d$ - $2\sigma_p$ 
  - Interpretation does not require knowledge of parton distributions.

$$A_{d2p} = \frac{\sigma_d^L - \sigma_d^R - 2(\sigma_p^L - \sigma_p^R)}{\sigma_d^L + \sigma_d^R - 2(\sigma_p^L + \sigma_p^R)}$$

$$= \left(\frac{G_F Q^2}{\pi \alpha 2 \sqrt{2}}\right) \left[-\frac{1}{2} + 2\sin^2(\theta_W)\right]$$

$$\times [1 + Y]$$

$$\approx -0.65 \times 10^{-5} Q^2 (1 + Y)$$

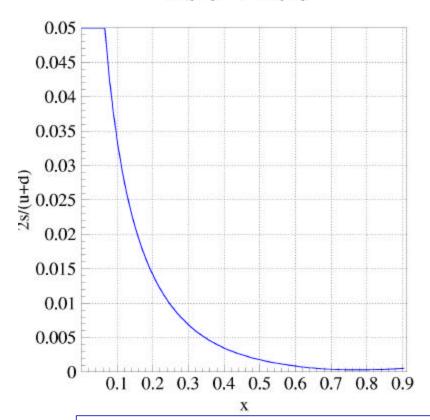
- Ratio of asymmetries: A<sub>p</sub>/A<sub>d</sub>
  - If  $C_{1a}$ 's are known, measures  $r(x) \frac{1}{4} \frac{d(x)}{u(x)}$  at large x.
  - Polarization cancels out.

$$\left(\frac{A_p}{A_d}\right) = \left(\frac{2C_{1u} - r(x)C_{1d}}{2C_{1u} - C_{1d}}\right) \left(\frac{5}{4 + r(x)}\right)$$
$$r(x) \approx d(x)/u(x)$$

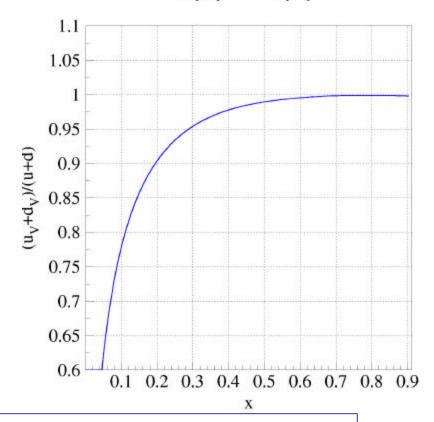
- s-quark distribution at low x: A<sub>p</sub>
  - Q<sup>2</sup> possibly not high enough at Jlab 11 GeV.

# $R_s(x)$ and $R_V(x)$

$$R_s(x) = \frac{2s(x)}{u(x) + d(x)} \xrightarrow{\text{large } \times} 0$$



$$R_v(x) = \frac{u_v(x) + d_v(x)}{u(x) + d(x)} \xrightarrow{\text{large} \times} 1$$



Uncertainties in PDF's are now known and would be factored into overall error budget.